BROKEN-RAY TOMOGRAPHY

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Zero Scattering Regime: Conventional X-Ray Tomography



Strong Scattering Regime: Diffuse Optical Tomography



- Many source-detector pairs
- Severely ill-posed IP
- Nonlinear IP

Mesoscopic Scattering Regime: Single-Scattering Tomography



SSOT And Other Modalities

	Linear IP	IP is MILDLY Ill-posed	Single projection	Reflection geometry	Nonionizing radiation	Quantitatrive images	Reconstruc- tion of scattering and absorption
X-ray CT	Y (if <i>E</i> =const)	Y	N	Ν	Ν	Y	Ν
DOT	Ν	Ν	Y	Y	Υ	?	Y (with time or freq.)
SSOT	Y	Y	Y	Y	Υ	Υ	Y (with two det. angles)

The Broken-Ray Integral Transform

(a) $\mu_s = \text{const}$ (and is known) $\int_{\text{SSR}} \mu_t[\mathbf{r}(\ell)] d\ell = \phi(\mathbf{r}_1, \hat{\mathbf{s}}_1; \mathbf{r}_2, \hat{\mathbf{s}}_2)$

(b)
$$\mu_s \neq \text{const}$$
 (and is unknown)

$$\int_{\text{SSR}} \mu_t[\mathbf{r}(\ell)] d\ell - \ln\left[\frac{\mu_s(\mathbf{R}_{12})}{\langle \mu_s \rangle}\right]$$

$$= \phi(\mathbf{r}_1, \hat{\mathbf{s}}_1; \mathbf{r}_2, \hat{\mathbf{s}}_2)$$

The measurable data function:

$$\phi(\mathbf{r}_1, \hat{\mathbf{s}}_1; \mathbf{r}_2, \hat{\mathbf{s}}_2) = -\ln\left[\frac{r_{12}\sin\theta_1\sin\theta_2}{\langle\mu_s\rangle A(\hat{\mathbf{s}}_1, \hat{\mathbf{s}}_2)} \frac{I_{\text{measured}}}{I_{\text{incident}}}\right]$$



Broken rays: You can go two ways



Outline:

- 1) Numerical test (40x40 rays per slice), full RTE forward solver, no inverse crime
- 2) Generalized filtered back-projection formula
- Inverse crime simulations based on this formula (many rays)

PART 1: Numerical simulations: Data from the RTE (no inverse crime)

- Forward model based on the RTE
- Isotropic scattering
- FULL ACCOUNT OF MULTIPLE SCATTERING
- BOUNDARY CONDITIONS SATISFIED EXACTLY
- 3D integral equation for density discretized on a rectangular grid
- Direct inversion of a well-posed square matrix
- Mathematical details on next page...

$$\begin{bmatrix} \hat{\mathbf{s}} \cdot \nabla + \mu_{a}(\mathbf{r}) + \mu_{s}(\mathbf{r}) \end{bmatrix} I(\mathbf{r}, \hat{\mathbf{s}}) = \mu_{s}(\mathbf{r}) \int \frac{1}{4\pi} I(\mathbf{r}, \hat{\mathbf{s}}') d^{2} \hat{\mathbf{s}}'$$

$$\mu_{t}(\mathbf{r}) \qquad A(\hat{\mathbf{s}}, \hat{\mathbf{s}}') = \frac{1}{4\pi} = \text{const}$$

$$u(\mathbf{r}) = \int I(\mathbf{r}, \hat{\mathbf{s}}) d^{2} \hat{\mathbf{s}}$$

$$u(\mathbf{r}) = u_{b}(\mathbf{r}) + \int g_{b}(\mathbf{r}, \mathbf{r}') \frac{\mu_{s}(\mathbf{r}')}{4\pi} u(\mathbf{r}') d^{3} r'$$

$$I(\mathbf{r}, \hat{\mathbf{s}}) = I_{b}(\mathbf{r}, \hat{\mathbf{s}}) + \int G_{b}(\mathbf{r}, \hat{\mathbf{s}}; \mathbf{r}', \hat{\mathbf{s}}') \frac{\mu_{s}(\mathbf{r}')}{4\pi} u(\mathbf{r}') d^{3} r' d^{2} \hat{\mathbf{s}}'$$

$$\frac{G_{b}(\mathbf{r}, \hat{\mathbf{s}}; \mathbf{r}', \hat{\mathbf{s}}') = g_{b}(\mathbf{r}, \mathbf{r}') \delta(\hat{\mathbf{u}}(\mathbf{r} - \mathbf{r}') - \hat{\mathbf{s}}') \delta(\hat{\mathbf{s}} - \hat{\mathbf{s}}')}{g_{b}(\mathbf{r}, \mathbf{r}')} = \int G_{b}(\mathbf{r}, \hat{\mathbf{s}}; \mathbf{r}', \hat{\mathbf{s}}') d^{2} \hat{\mathbf{s}} d^{2} \hat{\mathbf{s}}' =$$

$$= \frac{1}{|\mathbf{r} - \mathbf{r}'|^{2}} \exp \left[-\int_{0}^{|\mathbf{r} - \mathbf{r}'|} \mu_{t} (\mathbf{r}' + \ell \hat{\mathbf{u}}(\mathbf{r} - \mathbf{r}')) d\ell \right]$$







 $\mu_s = 0.16h^{-1}$ $\mu_s L_z = 6.4$



Simultaneous reconstruction of absorption and scattering

SCATTERING

Background: $\overline{\mu}_{s}L_{z} = 1.6$ Targets: $1.33\overline{\mu}_{s} \le \mu_{s} \le 2\overline{\mu}_{s}$ ABSORPTION Background: $\overline{\mu}_a = 0.1 \overline{\mu}_s$ Targets: $2\overline{\mu}_a \le \mu_a \le 5\overline{\mu}_a$



Inhomogeneities with stronger scattering



Stronger absorption (overall)

SCATTERING Background: $\overline{\mu}_{s}L_{z} = 1.6$

Targets: $2\overline{\mu}_s \le \mu_s \le 3\overline{\mu}_s$

ABSORPTION Background: $\overline{\mu}_a = \overline{\mu}_s$ Targets: $2\overline{\mu}_a \le \mu_a \le 5\overline{\mu}_a$



Same as before, but larger optical depth



PART 2: FBF derivation



$$\eta(\Delta, \ell) = \begin{cases} 0, \ \ell < L_1(\Delta) \\ \ell - L_1(\Delta), \ L_1(\Delta) < \ell < L_1(\Delta) + L_2(\Delta) \end{cases}$$
$$\zeta(\Delta, \ell) = \begin{cases} \ell, \ \ell < L_1(\Delta) \\ \ell, \ \ell < L_1(\Delta) \end{cases}$$
$$\zeta(\Delta, \ell) = \{L_1(\Delta) + [\ell - L_1(\Delta)]\cos(\beta), \ L_1(\Delta) < \ell < L_1(\Delta) + L_2(\Delta) \end{cases}$$

 $L_1(\Delta) = L - \Delta \operatorname{ctg}(\beta)$, $L_2(\Delta) = \Delta / \sin(\beta)$

The ray langth (within the medium): $L_1(\Delta) + L_2(\Delta)$

Fourier slice theorem

$$\tilde{\phi}(k,\Delta) = \int_{0}^{L_{1}(\Delta)+L_{2}(\Delta)} e^{ik\eta(\Delta,\ell)} \tilde{\mu}_{t}[k,\zeta(\Delta,\ell)]d\ell$$
$$\tilde{\phi}(k,\Delta) = \int_{0}^{L_{1}(\Delta)} \tilde{\mu}_{t}(k,\ell)d\ell + \frac{e^{ikL_{1}(\Delta)\operatorname{tg}\beta}}{\cos\beta} \int_{L_{1}(\Delta)}^{L} \tilde{\mu}_{t}(k,\ell)e^{-ik\ell\operatorname{tg}\beta}d\ell$$
Define new variables: $q = k\operatorname{tg}\beta$; $c = \cos\beta$

Define new variables:
$$q - k \lg \beta$$
, $c - \cos \beta$

$$f(z) = \tilde{\mu}_t (q \operatorname{ctg} \beta, z);$$

$$F(z) = \tilde{\phi} [q \operatorname{ctg} \beta, (L - z) \operatorname{tg} \beta]$$



Putting everything together...

$$\tilde{\mu}_{t}(k,z) = \sigma \left\{ H(k,z) - ik\sigma e^{-ik\sigma z} \int_{0}^{z} e^{i\sigma k\ell} H(k,\ell) dl \right\}$$
$$\sigma = \operatorname{ctg}\left(\frac{\beta}{2}\right);$$
$$H(k,z) = \left(\frac{\partial}{\partial\Delta} + ik\right) \tilde{\phi}(k,\Delta) \Big|_{\Delta = (L-z)\operatorname{tg}\beta}$$

The real-space inversion formula

$$\mu_{t}(y,z) = \int_{-\infty}^{\infty} \tilde{\mu}_{t}(k,z)e^{-iky}\frac{dk}{2\pi}$$

$$\mu_{t}(y,z) = \sigma \left\{ \left(\frac{\partial}{\partial \Delta} - \frac{\partial}{\partial y} \right) \phi(y,\Delta) + \tau \frac{\partial}{\partial y} \left[\phi(y + \sigma z, L \operatorname{tg} \beta) - \phi(y,\Delta) \right] - (1 + \tau) \frac{\partial}{\partial y} \int_{\Delta}^{L \operatorname{tg} \beta} \phi(y + \tau(\ell - \Delta), \ell) d\ell \right\} \Big|_{\Delta = (L-z) \operatorname{tg} \beta}$$

$$\sigma = \operatorname{ctg}\left(\frac{\beta}{2} \right); \quad \tau = \operatorname{ctg}\left(\frac{\beta}{2} \right) / \operatorname{tg}(\beta)$$

PART 3: Inverse-crime simulations

Reconstruction using the Fourier-space formula



Model *L/h*=40 *L/h*=400

$$\beta = \frac{\pi}{4}$$

Reconstruction of a Gaussian



$$\beta = \frac{\pi}{4}$$
$$\frac{h}{L} = \frac{0.3}{40}$$

 $\frac{w}{L} = \frac{3}{40} \quad \frac{w}{L} = \frac{5}{40} \quad \frac{w}{L} = \frac{7}{40} \quad \frac{w}{L} = \frac{10}{40}$

Profiles of the reconstruction of a Gaussian



Real-space formula (for Gaussians)



Simultaneous reconstruction of absorption and scattering



(b)
$$\mu_s \neq \text{const}$$
 (and is unknown)

$$\int_{\text{SSR}} \mu_t[\mathbf{r}(\ell)] d\ell - \ln\left[\frac{\mu_s(\mathbf{R}_{12})}{\langle \mu_s \rangle}\right]$$

$$= \phi(\mathbf{r}_1, \hat{\mathbf{s}}_1; \mathbf{r}_2, \hat{\mathbf{s}}_2)$$

$$\phi = \phi_1 - \phi_2$$

Fourier-space image reconstruction formula for two rays

 $\tilde{\mu}_{t}(k,z) = \frac{\sin\beta}{2} \left(ik - \frac{1}{ik} \frac{\partial^{2}}{\partial\Delta^{2}} \right) \tilde{\phi}(k,\Delta) \Big|_{\Lambda=0}$







Same as above but for $\frac{\overline{\mu}_s}{\overline{\mu}_a} = 1$

SUMMARY

- SSOT allows accurate quantitative reconstruction of the attenuation function.
- With additional measurements, scattering and absorption can be reconstructed separately
- Ill-posedness of the inverse problem is very mild.
- Tomographic imaging is feasible up to about six scattering lengths, with the noise-to-signal level of about 3% or less.

Preliminary Experiment







Figure 6: Experimental measurements of the specific intensity as a function of the exit position on the slab surface for intralipid concentrations 0.02% (a) and 0.04% (b).