Tomography of Highly Scattering Media with the Method of Rotated Reference Frames

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Motivation (The Forward Problem Perspective)

	Plane Wave Modes	The Weyl Expanson
The Helmholtz Equation $(\nabla^2 + k_0^2)u = 0$	$e^{i\mathbf{k}\cdot\mathbf{r}}$	$\frac{e^{ik_0r}}{r} = \frac{i}{2\pi} \int Q^{-1} e^{i(\mathbf{q}\cdot\mathbf{p}+Qz)} d^2 q$
k = const	$\mathbf{k} \cdot \mathbf{k} = k_0^2$	$Q = \sqrt{k_0^2 - q^2}$
The Diffusion Equation	$e^{-\mathbf{k}\cdot\mathbf{r}}$	$\frac{e^{-k_0 r}}{e^{-k_0 r}} = \frac{1}{1} \int O^{-1} e^{i\mathbf{q}\cdot\mathbf{p}-Qz} d^2 a$
$(-\nabla \cdot D\nabla + \alpha)u = 0$	$\mathbf{k} \cdot \mathbf{k} - k^2 - \alpha / D$	$r \frac{2\pi}{2\pi} \int \mathcal{Q} c u q$
$D, \alpha = \text{const}$	$\mathbf{K} \cdot \mathbf{K} = \kappa_0 = \alpha + D$	$Q = \sqrt{k_0^2 + q^2}$
RTE	Not known	Not known
$(\hat{\mathbf{s}} \cdot \nabla + \mu_t) I(\mathbf{r}, \hat{\mathbf{s}}) =$		
$= \mu_s \int A(\hat{\mathbf{s}}, \hat{\mathbf{s}}') I(\mathbf{r}, \hat{\mathbf{s}}')$		
$\mu_t, \mu_s = \text{const}$		

MOTIVATION (The Inverse Problems Perspective)



Given a data function $\phi(\mathbf{\rho}_s,\mathbf{\rho}_d)$ which is measured for multiple pairs $(\mathbf{\rho}_s, \mathbf{\rho}_d)$, find the absorption coefficient $\alpha(\mathbf{r})$ inside the slab

Linearized Integral Equation

$$\phi(\mathbf{\rho}_s,\mathbf{\rho}_d) = \int \Gamma(\mathbf{\rho}_s,\mathbf{\rho}_d;\mathbf{r})\delta\alpha(\mathbf{r})d^3r$$

$$\phi(\mathbf{\rho}_{s}, \mathbf{\rho}_{d}) = \frac{I(\mathbf{\rho}_{s}, z_{s}; \mathbf{\rho}_{d}, z_{d}) - I_{0}(\mathbf{\rho}_{s}, z_{s}; \mathbf{\rho}_{d}, z_{d})}{I_{0}(\mathbf{\rho}_{s}, z_{s}; \mathbf{\rho}_{d}, z_{d})}$$
(measurable data-function)

 $\Gamma(\mathbf{\rho}_{s}, \mathbf{\rho}_{d}) = G_{0}(\mathbf{\rho}_{s}, z_{s}; \mathbf{r})G_{0}(\mathbf{r}; \mathbf{\rho}_{d}, z_{d})$ (first Born approximation) $\alpha(\mathbf{r}) = \alpha_{0} + \delta\alpha(\mathbf{r})$

Analytical SVD approach: Making use of the translational invariance

$$\tilde{\phi} (\mathbf{q}_{s}, \mathbf{q}_{d}) = \int \phi (\mathbf{\rho}_{s}, \mathbf{\rho}_{d}) e^{i(\mathbf{q}_{s} \cdot \mathbf{\rho}_{s} + \mathbf{q}_{d} \cdot \mathbf{\rho}_{d})} d^{2} \rho_{s} d^{2} \rho_{d}$$

$$\mathbf{q}_{s} = \mathbf{q} / 2 + \mathbf{p} , \quad \mathbf{q}_{d} = \mathbf{q} / 2 - \mathbf{p};$$
Data function: $\psi(\mathbf{q}, \mathbf{p}) = \tilde{\phi} (\mathbf{q} / 2 + \mathbf{p}, \mathbf{q} / 2 - \mathbf{p})$

$$\psi(\mathbf{q}, \mathbf{p}) = \int_{0}^{L} g_{s} (\mathbf{q} / 2 + \mathbf{p}; z) g_{d} (\mathbf{q} / 2 - \mathbf{p}; z) \delta \tilde{\alpha} (\mathbf{q}; z) dz$$

$$\delta \tilde{\alpha} (\mathbf{q}; z) dz = \int \delta \alpha (\mathbf{\rho}, z) e^{i\mathbf{\rho} \cdot \mathbf{q}} d^{2} \rho$$

$$G_{0} (\mathbf{\rho}_{s}, z_{s}; \mathbf{\rho}, z) = \int \frac{d^{2} q}{(2\pi)^{2}} g_{s} (\mathbf{q}; z) e^{i\mathbf{q} \cdot (\mathbf{\rho} - \mathbf{\rho}_{s})}$$

$$G_{0} (\mathbf{\rho}, z; \mathbf{\rho}_{d}, z_{d}) = \int \frac{d^{2} q}{(2\pi)^{2}} g_{d} (\mathbf{q}; z) e^{i\mathbf{q} \cdot (\mathbf{\rho} - \mathbf{\rho}_{s})}$$



S.D.Konecky, G.Y.Panasyuk, K.Lee, V.Markel, A.G.Yodh and J.C.Schotland Imaging complex structures with diffuse light *Optics Express* 16(7), 5048-5060 (2008) In the case of RTE, we need the plane-wave decomposition of the Greens function, of the form

$$G_0(\mathbf{r},\hat{\mathbf{s}};\mathbf{r}',\hat{\mathbf{s}}') = \int \frac{\mathrm{d}^2 q}{\left(2\pi\right)^2} g(\mathbf{q};z,\hat{\mathbf{s}};z',\hat{\mathbf{s}}') e^{i\mathbf{q}\cdot(\boldsymbol{\rho}-\boldsymbol{\rho}')}$$

and the integral kernel

$$\Gamma(\mathbf{q},\mathbf{p};z) = \int g(\mathbf{q} / 2 + \mathbf{p};z_s, \hat{\mathbf{z}};z, \hat{\mathbf{s}})g(\mathbf{q} / 2 - \mathbf{p};z, \hat{\mathbf{s}};z_d, \hat{\mathbf{z}})d^2s$$

We then get the integral equations of the form

$$\psi(\mathbf{q},\mathbf{p}) = \int_{z_s}^{z_d} \Gamma(\mathbf{q},\mathbf{p};z) \delta \tilde{\alpha}(\mathbf{q};z) dz$$

THE METHOD

Rotated Reference Frames

The usual sperical harmonics are defined in the laboratory reference frame. Then θ and φ are the polar angles of the unit vector \hat{s} in that frame.

THE MAIN IDEA:

For each value of the Fourier variable **k**, use spherical harmonics defined in a reference frame whose z-axis is aligned with the direction of **k**.

We call such frames "rotated". Spherical harmonics defined in the rotaded frame are denoted by $Y(\hat{\mathbf{s}}; \hat{\mathbf{k}})$.

Rotation of the Laboratory Frame (x,y,z).



Advantages of the Method

- Numerical problem is reduced to diagonalization of a set of tridiagonal matrices
- Once the eigenvalues and eigenvectors are computed, the Green's function, the plane-wave modes and the Weyl expansion for the GF can be obtained in analytical form
- The Weyl expansion can also be obtained in a slab with appropriate boundary conditions

RESULTS: SIMULATIONS



A set of 5 point absorbers in an L=6*l** slab The field of view is 16*l**



A bar target in the center of the same slab



RTE

Diffusion approximation

RESULTS: SIMULATIONS

Two thin vertical wires in a 1cm thick slab filled with intralipid solution (looks like milk)





Diffusion approximation

Publications:

1. V.A.Markel, "Modified spherical harmonics method for solving the radiative transport equation," *Waves in Random Media* **14**(1), L13-L19 (2004).

2. G.Y.Panasyuk, J.C.Schotland, and V.A.Markel, "Radiative transport equation in rotated reference frames," *Journal of Physics A*, **39**(1), 115-137 (2006).

3. J.C.Schotland and V.A.Markel, "Fourier-Laplace structure of the inverse scattering problem for the radiative transport equation," Inverse Problems and Imaging 1(1), 181-188 (2007).

Available on the web at http://whale.seas.upenn.edu/vmarkel/papers.html

CONCLUSIONS

- The method of rotated reference frames can be used in optical tomography of mesoscopic samples
- The images are of superior quality compared to those obtained by using the diffusion approximation
- The quality of images can be comparable to that in X-ray tomography because the RTE retains some information about ballistic rays, single-scattered rays, etc.